# Epitaxial Growth of LaB<sub>6</sub> Thin Films on Ultrasmooth Sapphire Substrate with Epitaxial SrB<sub>6</sub> Buffer Layer

Yushi KATO<sup>\*1</sup>, Naoki SHIRAISHI<sup>\*1</sup>, Satoru KANEKO<sup>\*2</sup>, Nobuo TSUCHIMINE<sup>\*3</sup>, Susumu KOBAYASHI<sup>\*3</sup> and Mamoru YOSHIMOTO<sup>\*1, \*4</sup>

<sup>\*1</sup> Department of Innovative & Engineered Materials, Tokyo Institute of Technology, 4259-J2-46, Nagatsuta-cho, Midori-ku, Yokohama 226-8503, Japan

E-mail:kato.y.am@m.titech.ac.jp

<sup>\*2</sup> Kanagawa Industrial Technology Research Institute, 705-1 Shimo-Imaizumi, Ebina, Kanagawa 243-0435, Japan

<sup>\*3</sup> TOSHIMA Manufacturing Company Limited, 1414, Shimonomoto, Higashimatsuyama-shi, Saitama 355-0036, Japan

<sup>\*4</sup> Patent Attorney, Tokyo Institute of Technology, 4259-J2-46, Nagatsuta-cho, Midori-ku, Yokohama 226-8503, Japan

We fabricated epitaxial LaB<sub>6</sub> (100) thin films on ultrasmooth sapphire ( $\alpha$ -Al<sub>2</sub>O<sub>3</sub> single crystal) (0001) substrate with epitaxial SrB<sub>6</sub> buffer layer by laser molecular beam epitaxy. Reflection highenergy electron diffraction and X-ray diffraction measurements indicate the heteroepitaxial structure of LaB<sub>6</sub> (100)/SrB<sub>6</sub> (100)/sapphire (0001) with three domains of epitaxial relationship. The prepared films exhibit atomically stepwise surface morphology, similar to that for the substrate, with 0.2-nmhigh atomic steps and ~70-nm-wide terraces. LaB<sub>6</sub> epitaxial thin films show metallic behavior, with almost constant resistivity ( $1.0 \times 10^{-3} \Omega$  cm) in the temperature range 10–300 K, while epitaxial SrB<sub>6</sub> buffer thin films show semiconducting behavior, with a resistivity of 4.8  $\Omega$  cm at room temperature.

Keywords: hexaboride, LaB<sub>6</sub>, SrB<sub>6</sub>, laser MBE, epitaxial thin film, electrical properties

## 1. Introduction

Alkaline hexaborides, such as  $SrB_6$ ,  $CaB_6$ , and  $BaB_6$ , and rare earth hexaborides, such as  $LaB_6$ ,  $EuB_6$ , and  $CeB_6$ , have attractive physical properties. Their structure is a CsCl-type simple cubic arrangement composed of  $B_6$  octahedra and metal ions [1-4]. Interestingly, their electrical properties vary widely, from semiconducting [5] to superconducting [6], upon doping divalent or trivalent metallic ions into the crystal structure.

LaB<sub>6</sub> has a high melting temperature, excellent thermal stability, high hardness, and excellent chemical stability [1-4,7-9]. It is often used in thermionic electron sources for transmission electron microscopy, scanning electron microscopy, and flat panel displays, for which it offers high brightness and long service life because of its extremely small work function (~2.4 eV), high current and voltage capability, and low vapor pressure at high temperature. Its electric conduction reportedly shows metallic properties [1,7,8].

Investigations into the properties of  $LaB_6$  have been conducted mostly on bulk samples [1-4,7-9]. The properties of  $LaB_6$  thin films have not been well explored. There have been a few attempts to prepare  $LaB_6$  nanowire by chemical vapor deposition [10] and  $LaB_6$  polycrystalline thin films by sputtering [11,12], e-beam evaporation [13], and pulsed laser deposition (PLD) [14-17]. However, there are few reports on fabrication of epitaxial  $LaB_6$  thin films [18], although heteroepitaxial growth of hexaboride thin films is of great interest for their applicability in thermoelectric conversion materials [19] and high-performance electronic devices.

We previously investigated low-temperature epitaxial growth of functional ceramic thin films of oxides [20-22], nitrides [23,24], and diamond [25] by PLD or laser molecular beam epitaxy (MBE) (that is, pulsed laser deposition in ultrahigh vacuum) on ultrasmooth sapphire ( $\alpha$ -Al<sub>2</sub>O<sub>3</sub> single crystal) (0001) substrate. The smooth substrate has 0.2-nm-high atomic steps and 60–80-nm-wide atomically flat terraces [26]. Laser MBE has been used extensively to grow epitaxial thin films of high-melting-point ceramics at relatively low substrate temperatures [20-25].

Recently, we fabricated highly (100)-oriented LaB<sub>6</sub> polycrystalline thin films and epitaxial SrB<sub>6</sub> thin films on the ultrasmooth sapphire (0001) substrate by laser MBE [17]. In this study, we grew epitaxial LaB<sub>6</sub> thin films on ultrasmooth sapphire (0001) substrate with epitaxial SrB<sub>6</sub> buffer layer by laser MBE and characterized the film's morphological, crystallographic, and electrical properties.

# 2. Experimental details

Fabrication of SrB<sub>6</sub> and LaB<sub>6</sub> thin films was carried out by laser MBE ( $\sim 10^{-7}$  Pa base pressure) on ultrasmooth sapphire (0001) substrate. The substrate was obtained by annealing commercially available mirror-polished  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> single crystal wafers at 1000 °C for 3 h in air [26]. The step height was uniformly ~0.2 nm, and the terrace width was 60–80 nm. A pulsed KrF excimer laser (248 nm wavelength, 20 ns pulse duration, 0.5 mm<sup>2</sup> beam spot size) was focused on a sintered ceramic target of  $SrB_6$  and  $LaB_6$  (both 99.5% purity), with energy densities of 1.5 and 3.0 J/cm<sup>2</sup> and frequencies of 3 and 5 Hz, respectively. The distance between the substrate and the target was 5 cm. The ablated particles collided with the substrate heated at 800 °C to form the films. The elapsed time of deposition for  $SrB_6$  and  $LaB_6$  was 30 and 60 min, respectively.

Crystallinity and surface atomic structure during film deposition were characterized by in situ reflection highenergy electron diffraction (RHEED). Crystallographic structure was characterized by ex situ X-ray diffraction (XRD; MXP-M18, Bruker AXS) using CuK $\alpha$  line ( $\lambda = 0.1542$  nm). The film surface morphology was observed in air by atomic force microscopy (AFM; SPI-3700, Seiko Instruments). The film surface work function  $\varphi$  was determined using a Kelvin probe (KP) in air. The film resistivity was measured by the conventional four-probe method in the temperature range 10–300 K.

#### 3. Results and discussion

First, we examined the epitaxial growth of  $LaB_6$  films deposited directly on sapphire substrate. Fig. 1 shows the XRD pattern and RHEED image for a  $LaB_6$  thin film (30 nm) on sapphire substrate without epitaxial SrB<sub>6</sub> buffer layer. From the XRD pattern and the ring RHEED image,  $LaB_6$  thin film was not epitaxial, but the (100) highly oriented polycrystalline film with the in-plane random configuration. We found that the lattice mismatch (12.7%) between sapphire and  $LaB_6$  is very large to carry out successful epitaxial growth of  $LaB_6$  films on sapphire (0001) substrate.

Next, we considered the possibilities of SrB<sub>6</sub> buffer layer to reduce lattice mismatch between LaB<sub>6</sub> and sapphire (0001), because the lattice mismatch between  $LaB_6$ and SrB<sub>6</sub> is just 0.9%. Fig. 2 shows the XRD pattern and RHEED image for a SrB<sub>6</sub> thin film (15 nm) on sapphire substrate. The streaky RHEED image and evident in-plane anisotropy indicate epitaxial film growth. From the XRD pattern,  $SrB_6$  (100) and (300) peaks are observed, and the  $SrB_6$  (200) peak overlaps with the sapphire (0006) peak. The root-mean-square (RMS) roughness of the film on the substrate is 0.10 nm. The film exhibits atomically stepwise surface morphology, similar to that for the substrate, with 0.2-nm-high atomic steps and ~70-nm-wide terraces. Thus, epitaxial SrB<sub>6</sub> (100) thin films can be fabricated on ultrasmooth sapphire substrate and are suitable for use as a buffer layer.

To determine the influence of an ultrathin epitaxial  $SrB_6$ buffer layer, we grew  $LaB_6$  thin films on ultrasmooth sapphire substrate, with and without epitaxial  $SrB_6$  buffer layer, under the same conditions. Fig. 3 shows the XRD pattern and RHEED image for a  $LaB_6$  thin film (30 nm) on sapphire substrate with epitaxial  $SrB_6$  buffer layer (5 nm). The RHEED image is streaky, and the image intensity and pattern change with incident beam direction. The streaky image and evident in-plane anisotropy indicate epitaxial film growth. From the XRD pattern,  $LaB_6$  (100) and (200) planes are observed. Although slight reflection peak of  $LaB_6$  (111) was observed, it was found that  $LaB_6$  thin films with single phase were grown preferentially compared to the growth without  $SrB_6$  buffer, as shown in insets of the Fig. 3. Thus, epitaxial  $LaB_6$  (100) thin films can be fabricated on ultrasmooth sapphire substrate with epitaxial  $SrB_6$ buffer layer, and the buffer layer promotes epitaxial growth of the films.



Fig. 1 XRD pattern of LaB<sub>6</sub> thin film on sapphire substrate without  $SrB_6$  buffer layer. Inset: RHEED image viewed along the sapphire [10-10] axis.



**Fig. 2** XRD pattern of  $SrB_6$  thin film on sapphire substrate. Inset: RHEED image viewed along the sapphire [10-10] axis.



**Fig. 3** XRD pattern of LaB<sub>6</sub> thin film on sapphire substrate with SrB<sub>6</sub> buffer layer. Insets: magnified ( $\times$ 35) view of XRD pattern in the 2 theta range 25-40 deg. and RHEED image viewed along the sapphire [10-10] axis.

We investigated the crystal structure and in-plane orientation of the obtained LaB<sub>6</sub> epitaxial thin film by rotary  $\varphi$ scan XRD using a four-circle X-ray diffractometer. Fig. 4 shows the azimuthal  $\varphi$ -scan XRD result for LaB<sub>6</sub> (101) peaks. Peaks (circle, triangle, square) are observed at every 60° phi rotation. Small peaks (dotted circle, dotted triangle, dotted square) are similarly observed. For a simple cubic LaB<sub>6</sub>/SrB<sub>6</sub> layer grown on sapphire substrate, the small peaks will be more clearer, indicating that the LaB<sub>6</sub> thin film and SrB<sub>6</sub> buffer layer grow as a tetragonal-like cubic, not simple cubic, structure since the in-plane lattice misfit over the two materials (SrB<sub>6</sub> and sapphire) is -11.9 and 1.8% in the [1-100] and [11-20] directions, respectively, of sapphire.

Thus, three domains of  $LaB_6/SrB_6$  layer on sapphire substrate were grown with in-plane two-folded symmetry. The critical thickness between  $SrB_6$  and sapphire was estimated theoretically by the People–Bean formula [27]

$$h_c \approx (\frac{1-\nu}{1+\nu})(\frac{1}{16\pi\sqrt{2}})(\frac{b^2}{a_{epi}})[(\frac{1}{f^2})\ln(\frac{h_c}{b})],\tag{1}$$

where  $h_c$  is the critical thickness, v is the Poisson ratio, b is the slip distance,  $a_{epi}$  is the lattice constant for the SrB<sub>6</sub> epitaxial layer, and f is the lattice mismatch between SrB<sub>6</sub> and sapphire. The critical thickness (~40 nm) in the [11-20] direction of sapphire can be calculated, while that in the [1-100] direction cannot since the misfit is too large. This implies that strain relaxation begins immediately after film deposition in the [1-100] direction of sapphire; in contrast, it will not begins until 40 nm film thickness in the [11-20] direction of sapphire [28]. In fact, the crystalline films of the LaB<sub>6</sub>/SrB<sub>6</sub>/sapphire and SrB<sub>6</sub>/sapphire systems turn from epitaxial to polycrystalline with increasing film thickness. XRD and RHEED measurements indicate the heteroepitaxial structure of LaB<sub>6</sub> (100)/SrB<sub>6</sub> (100)/sapphire (0001) with three domains of epitaxial relationship: LaB<sub>6</sub>  $[001]//SrB_6$  [001]//sapphire [1-100],  $LaB_6$   $[001]//SrB_6$ [001]//sapphire [10-10], and LaB<sub>6</sub>  $[001]//SrB_6$ [001]//sapphire [01-10].



**Fig. 4**  $\varphi$ -scan XRD pattern of LaB<sub>6</sub> thin film on sapphire (0001) substrate with SrB<sub>6</sub> (100) buffer layer.



Fig. 5 AFM surface image of  $LaB_6$  thin film on sapphire substrate with  $SrB_6$  buffer layer.

Fig. 5 shows the AFM surface image of the prepared LaB<sub>6</sub> thin film on sapphire substrate with epitaxial SrB<sub>6</sub> buffer layer. The average diameter of the observed grainlike islands agrees well with that calculated by Sherrer's equation ( $\sim$ 20 nm). The film surface is atomically flat (RMS roughness 0.28 nm) and exhibits straight atomic steps derived from the surface of the substrate.

From the earliest considerations of the electronic structure of metal hexaborides, it has been predicted that those formed from trivalent metals would be metallic electrical conductors and those formed from divalent metals would be semiconductors [7,29]. Fig. 6 shows the temperature dependence of resistivity for a SrB<sub>6</sub> thin film on sapphire substrate; the inset shows a magnified view of the curve in the temperature range 250–300 K. The film shows semiconducting behavior, and the resistivity is about 5  $\Omega$  cm at room temperature. Two different regions of activated behavior can be identified [5]. Low-temperature and hightemperature regions were fit independently to an activated form by the Arrhenius equation

$$\rho(T) = \rho_0 \exp(\Delta E / k_B T), \qquad (2)$$

where  $\Delta E$  is the activation energy,  $\rho_0$  is the resistivity at large temperature *T*. The fit results are in excellent agreement with a semiconducting model in which the hightemperature fit gives a gap of ~38.2 meV and the lowtemperature fit gives a much smaller gap of ~0.5 meV. These gaps are about 10 times larger than for a singlecrystal bulk sample (5.8 meV in the high-temperature range and 0.09 meV in the low-temperature range) [5].

Fig. 7 shows the temperature dependence of resistivity for a LaB<sub>6</sub> thin film on sapphire substrate with epitaxial SrB<sub>6</sub> buffer layer. The film shows metallic behavior, and the resistivity is almost constant  $(1.0 \times 10^{-3} \Omega \text{ cm})$  in the temperature range 10–300 K. The room-temperature resistivities for the epitaxial LaB<sub>6</sub> and SrB<sub>6</sub> films are about 100 times larger than those reported for single-crystal bulk LaB<sub>6</sub> and SrB<sub>6</sub> [1,5].

In general, the difference in resistivity for bulk crystal and thin film is considered to be caused by various effects of impurity, deficiency, internal stress, and compositional nonstoichiometry that exist in thin films. For alkaline earth hexaborides such as  $SrB_6$ , the valence-band maximum and conduction-band minimum are sensitive to the intra B-octahedron bond length; therefore, the band gap depends on the lattice parameter [30]. Thus, considering that the lattice parameter of the films ( $a_0 = 0.4255$  nm) is larger than that of bulk  $SrB_6$  ( $a_0 = 0.4195$  nm), resistivity is affected by the lattice parameter.

On the other hand, from the band calculation for LaB<sub>6</sub>, the Fermi surface consists of strongly hybridized orbitals formed by hybridization between boron *sp* orbitals and lanthanum *d* orbitals, so boron vacancies can seriously disturb electron conduction paths [13]. To determine the influence of electron conduction paths, we calculated the mean free path of an electron carrier in our sample and compared it with that for the bulk sample reported previously [8,13]. The mean free path *l* of an electron carrier in our sample was estimated by Matthiessen's rule [13]

$$l = (r.r.r.-1)[\rho l]_{Bulk} / \rho_{(T=300K)}, \qquad (3)$$

where  $[\rho l]_{Bulk}$  is the bulk value of  $\rho l$  (2.1 × 10<sup>-11</sup>  $\Omega$  cm<sup>2</sup> for LaB<sub>6</sub> [8]), *r.r.* is the residual resistance ratio, and  $\rho_{(T=300 K)}$  is the bulk resistivity at 300 K (8.9 × 10<sup>-6</sup>  $\Omega$  cm for LaB<sub>6</sub>). Thus, the mean free path *l* for the film is about 0.8 nm, much smaller than that for the bulk material (~6.98 µm [8,13]). Since the grain size *d* (~20 nm) from the XRD and AFM results is larger than the mean free path *l* estimated for this sample, the effect of grain-boundary scattering on electron conduction is considered to be smaller than the effect of inter-grain scattering, including electron–electron, electron–phonon, and impurity/defect scattering. The relationship between electrical properties and composition in the thin film system is not clear at present and requires further investigation.

When examining LaB<sub>6</sub> thin films with a view toward their application as cathode materials and electronic devices, an important property is the work function  $\varphi$ . For reference, we measured  $\varphi$  for gold thin film, which is known to be good standard [31,32] and whose value we assumed to be 5.28 eV. From KP results, we estimated  $\varphi$ for bulk LaB<sub>6</sub> and LaB<sub>6</sub> thin film to be 4.53 and 4.54 eV, respectively.



Fig. 6 Temperature dependence of resistivity for  $SrB_6$  thin film on sapphire substrate. Inset: magnified view in the temperature range 250–300 K.



Fig. 7 Temperature dependence of resistivity for  $LaB_6$  thin film on sapphire substrate with  $SrB_6$  buffer layer.

These values do not agree with the previous reported value for the (001) clean surface of single-crystal bulk LaB<sub>6</sub> (2.4-2.6 eV) [2,3]. Since our measurements were performed in air, the surface of the sample should be covered with adsorbates such as oxygen and hydrogen, which may affect measurement. Further studies are thus necessary to determine the surface properties of hexaboride thin films.

## 4. Conclusions

We fabricated epitaxial LaB<sub>6</sub> (100) thin films on ultrasmooth sapphire ( $\alpha$ -Al<sub>2</sub>O<sub>3</sub> single crystal) (0001) substrate with epitaxial SrB<sub>6</sub> buffer layer. RHEED and XRD measurements indicate the heteroepitaxial structure of LaB<sub>6</sub> (100)/SrB<sub>6</sub> (100)/sapphire (0001) with three domains of epitaxial relationship: LaB<sub>6</sub> [001]//SrB<sub>6</sub> [001]//sapphire [1-100], LaB<sub>6</sub> [001]//SrB<sub>6</sub> [001]//sapphire [10-10], and LaB<sub>6</sub> [001]//SrB<sub>6</sub> [001]//sapphire [01-10]. The films exhibit atomically stepwise surface morphology, similar to that of the substrate, with 0.2-nm-high atomic steps and ~70-nmwide terraces. They are epitaxial in nature and show metallic behavior, with almost constant resistivity ( $1.0 \times 10^{-3} \Omega$ cm) in the temperature range 10–300 K. In contrast, the epitaxial SrB<sub>6</sub> thin films show semiconducting behavior, with a resistivity of 4.8  $\Omega$  cm at room temperature.

#### Acknowledgments

This work was supported in part by the Ministry of Education, Culture, Sports, Science and Technology of Japan, the National Institute of Advanced Industrial Science and Technology of Japan, the New Energy and Industrial Technology Development Organization of Japan, and the Regional Innovation Creation R&D Program of the Ministry of Economy, Trade and Industry of Japan.

### References

- D. Mandrus, B. C. Sales and R. Jin: Phys. Rev. B, 64, (2001) 012302.
- [2] M. Aono, R. Nishitani, C. Oshima, T. Tanaka, E. Bannai and S. Kawai: J. Appl. Phys., 50, (1979) 4802.
- [3] M. Aono, C. Oshima, T. Tanaka, E. Bannai and S. Kawai: J. Appl. Phys., 49, (1978) 2761.
- [4] C. -H. Chen, T. Aizawa, N. Iyi, A. Sato and S. Otani: Journal of Alloys and Compounds, 366, (2004) L6.

- [5] H. R. Ott, M. Chernikov, E. Felder, L. Degiorgi, E. G. Moshopoulou, J. L. Sarrao and Z. Fisk: Z. Phys. B, 102, (1997) 337.
- [6] R. J. Sobczak and M. J. Sienko: J. Less-Com. Metals, 67, (1979) 167.
- [7] A. Hasegawa and A. Yanase: J. Phys. F: Metal Phys., 7, (1977) 1245.
- [8] T. Tanaka, T. Akahane, E. Bannai, S. Kawai, N. Tsuda and Y. Ishizawa: J. Phys. C: Solid State Phys., 9, (1976) 1235.
- [9] M. Zhang, L. Yuan, X. Wang, H. Fan, X. Wang, X. Wu, H. Wang and Y. Qian: J. Solid State Chem., 181, (2008) 294.
- [10] H. Zhang, Q. Zhang, J. Tang and L. -C. Qin: J. Am. Chem. Soc., 127, (2005) 2862.
- [11] C. Mitterer: J. Solid State Chem., 133, (1997) 279.
- [12] T. Nakano, S. Baba, A. Kobayashi, A. Kinbara, T. Kajiwara and K. Watanabe: J. Vac. Sci. Technol. A, 9, (1991) 547.
- [13] I. Terasaki, S. Uchida, S. Tajima, K. Uchinokura and S. Tanaka: J. Phys. Soc. Jpn., 59, (1990) 1017.
- [14] V. Craciun and D. Craciun: Appl. Surf. Sci., 247, (2005) 384.
- [15] D. J. Late, K. S. Date, M. A. More, P. Misra, B. N. Singh, L. M. Kukreja, C. V. Dharmadhikari and D. S. Joag: Nanotechnology, 19, (2008) 265605.
- [16] D. J. Late, M. A. More, P. Misra, B. N. Singh, L. M. Kukreja and D. S. Joag: Ultramicroscopy, 107, (2007) 825.
- [17] Y. Kato, Y. Ono, Y. Akita, M. Hosaka, N. Shiraishi, N. Tsuchimine, S. Kobayashi and M. Yoshimoto: Mat. Res. Soc. Symp. Proc., 1148, (2009) PP12-02.
- [18] S. Muranaka and S. Kawai: J. Appl. Phys. Jpn., 15, (1976) 587.

- [19] M. Takeda, T. Fukuda, F. Domingo and T. Miura: J. Solid State Chem., 177, (2004) 471.
- [20] J. Tashiro, A. Sasaki, S. Akiba, S. Satoh, T. Watanabe, H. Funakubo and M. Yoshimoto: Thin Solid Films, 415, (2002) 272.
- [21] T. Maeda, M. Yoshimoto, T. Ohnishi, G. H. Lee and H. Koinuma: J. Cryst. Growth, 177, (1997) 95.
- [22] A. Sasaki, W. Hara, A. Matsuda, N. Tateda, S. Otaka, S. Akiba, K. Saito, T. Yodo and M. Yoshimoto: Appl. Phys. Lett., 86, (2005) 231911.
- [23] W. Hara, J. Liu, A. Sasaki, S. Otaka, N. Tateda, K. Saito and M. Yoshimoto: Thin Solid Films, 516, (2008) 2889.
- [24] A. Sasaki, J. Liu, W. Hara, S. Akiba, K. Saito, T. Yodo and M. Yoshimoto: J. Mater. Res., 19, (2004) 2725.
- [25] M. Yoshimoto, K. Yoshida, H. Maruta, Y. Hishitani, H. Koinuma, S. Nishio, M. Kakihana and T. Tachibana: Nature, 399, (1999) 340.
- [26] M. Yoshimoto, T. Maeda, T. Ohnishi, O. Ishiyama, M. Shinohara, M. Kubo, R. Miura, A. Miyamoto and H. Koinuma: Appl. Phys. Lett., 67, (1995) 2615.
- [27] R. People and J. C. Bean: Appl. Phys. Lett., 47, (1985) 322.
- [28] J. Zhu, X. H. Wei, Y. Zhang and Y. R. Li: J. Appl. Phys., 100, (2006) 104106.
- [29] H. C. Longuet-Higgins and M. D. V. Roberts: Proc. R. Soc. London, Ser. A, 224, (1954) 336.
- [30] B. Lee and L. -W. Wang: Appl. Phys. Lett., 87, (2005) 262509.
- [31] R. Fujii, Y. Gotoh, M. Y. Liao, H. Tsuji and J. Ishikawa: Vacuum, 80, (2006) 832.
- [32] N. A. Surplice and R. J. D'Arcy: J. Phys. E, 3, (1970) 477.

(Received: July 9, 2009, Accepted: November 25, 2009)