

Laser Melting of Nanoparticulate Transparent Conductive Oxide Thin Films

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Layers of ZnO nanoparticles with thicknesses of about 40 nm were prepared on silicon substrates. The layers were irradiated by single pulses of a 248 nm excimer laser, which proved suitable for consolidation and significant densification of the particle layers under ambient conditions. Experiments as well as simulations have confirmed that the application of a SiO₂ particle layer between the substrate and the ZnO particle layer can be used to hamper heat transfer towards the substrate. Thus the ZnO layer can be thermally insulated from the substrate while heating the ZnO up to its extremely high melting point. Consequently, such a layer stack could enable the application of consolidated particle layers on temperature-sensitive carrier substrates such as polymer foils which are to be used in low-cost mass production of devices like displays or solar cells.

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1. Introduction

Transparent conductive oxides, such as tin-doped indium oxide (ITO) or doped ZnO, exhibit a high electrical conductivity while at the same time offering transparency in the visible spectral range. Due to these properties they are suitable for a wide range of applications in optoelectronic devices, such as LC-displays, touchscreens or solar cells [1]. In most commercial products, the required conducting films are created by vapor deposition. However, for the low cost mass production of such devices, avoidance of the costly vacuum deposition step would be beneficial. This can be achieved by solution processing of the material by various printing techniques. A common method is the deposition of thin films out of nanoparticle dispersions [2, 3]. Usually the layer properties, e.g. its electrical conductivity can be enhanced by a subsequent annealing step, which can induce melting or sintering of the particles. By such a consolidating treatment the charge carrier transport through the film is greatly facilitated by enlargement of the particle-particle interfaces [4,5]. However, if low cost substrates like polymer foils are to be used, the processing temperatures need to be kept low in order to protect these substrates from thermal damage. To overcome this problem, we propose UV laser treatment to restrict energy deposition to the surface of the particle layer [6]. Furthermore, we show both by experiments and simulations that a layer of SiO₂ particles sandwiched between the substrate and the ZnO particle layer can be

used as a thermally insulating buffer layer which prevents heat transfer towards the substrate and thus also contributes to the protection of the substrate. We chose zinc oxide nanoparticles as an exemplary metal oxide nanoparticle system since it is a well-studied semiconducting material [7] which can be deposited from a solution via nanoparticle dispersions.

2. Simulation

To simulate the temperature distribution and the temperature course within a nanoparticle layer during and after laser treatment, finite element simulations in COMSOL 3.5 were carried out. The effect of a SiO₂ buffer layer as well as the influence of different parameters such as pulse length and the gas surrounding the particle layer were investigated.

Stacks of a 40 nm thick ZnO particle layer consisting of ZnO particles with a diameter of about 25 nm followed by a buffer layer of a 220 nm thick SiO₂ layer consisting of SiO₂ particles with a diameter of 50 nm on a Si substrate (from now on referred to as “stack A”) and of a 40 nm thick ZnO particle layer with a particle diameter of 25 nm on a Si substrate (“stack B”) were considered in the simulation. The described model setup can be seen in Fig 1. The particular values for the layer thicknesses and particle sizes were taken to match the experimental part. For the energy input single pulses of a 248 nm Excimer laser with a pulse length of 30 ns were used. The wavelength of 248 nm cor-

responds to a photon energy of 5 eV which exceeds the band gap of ZnO of about 3.4 eV [7]. Thus, it is very well absorbed in ZnO due to interband transitions leading to a very shallow penetration depth of about 50 nm in bulk ZnO [8]. In the simulations the laser pulse was modeled as rectangular in time, which is in good agreement with its actual temporal shape. Absorption of laser light in Si and ZnO was calculated according to Beer-Lambert law. Reflection of the laser light at boundaries of different materials was estimated by the Fresnel equations, while the optical properties of the different particle layers were averaged according to Bruggemann's model [9]. It should be mentioned here that the absorbed laser energy is assumed to be an ideal heat source neglecting all other effects like radiative transfers. The laser fluence was set to 245 mJ/cm² which is the fluence needed to achieve good ZnO particle consolidation in the samples without the SiO₂ buffer layer (see experimental part, Fig. 6 d). The dissipation of heat in the particle layer was calculated by the heat transfer equation, also thermal radiation of the surfaces was taken into account. The temperature-dependent bulk thermal properties of ZnO, SiO₂, Si and air as well as the optical properties of ZnO and Si were taken from the literature [8, 10-14]. An appropriate heat sink for the incident laser energy was integrated by extending the lateral dimensions of the air side and the Si substrate side of the simulation area to 10 μm and 100 μm, respectively, resulting in constant temperatures at the limits of the simulation area during the whole simulated time frame of 1 μs. The transverse boundaries of the 150 nm wide simulation area were assumed to be thermally insulating resulting in periodic boundary conditions which should be reasonable as the laser spot is wider by several orders of magnitude. The calculation with the above parameters results in maximum temperatures of about 1200 °C for stack B, a temperature much lower than the melting point of ZnO of about 1975 °C [15]. Thus, according to the simulation the laser pulse would not have led to particle melting and consolidation. The reason for this underestimation of temperature is that the thermal conductivity of the used material was not represented correctly by the values of thermal conductivity for the bulk material. Since the used particles are only a few tens of nanometers in diameter, thermal conductivity i.e. phonon spread in the material is significantly reduced as the spatial dimension of the phonons is comparable to the particle size and the contact area between the particles. Consequently, this was taken into account by applying thermal conductivities measured on a comparable material [16]. This means that the

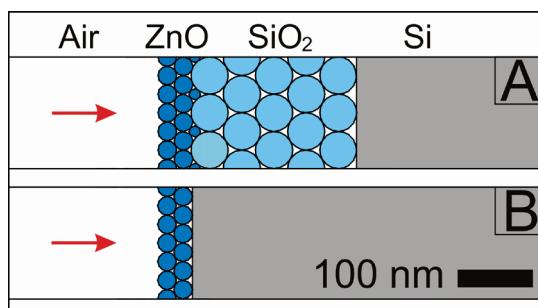


Fig. 1 Model setup used in the simulation (compare Fig. 6 (a) and Fig. 7(a)). The red arrows indicate the direction of the incident laser pulse

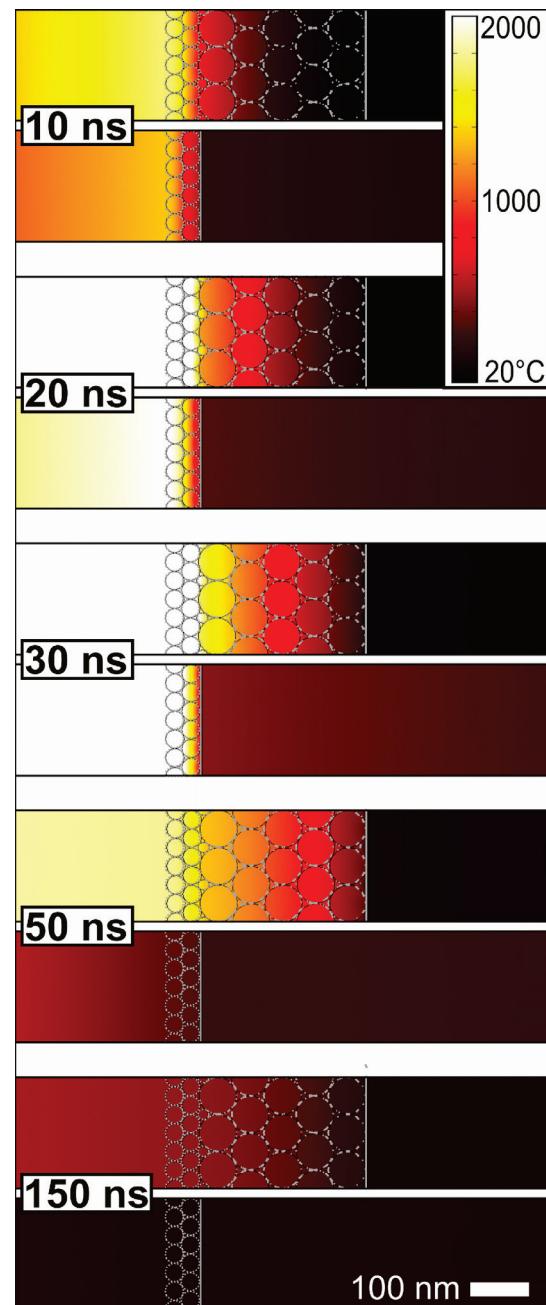


Fig. 2 Temperature distributions in the particle layers (see Fig. 1) at five exemplary time steps. Room temperature appears in black, while for clarity temperatures of 2000 °C and above are shown in white.

special properties of the material due to its nanoscale size are changed in our favor. Smaller particles should reduce the thermal conductivity even further while at the same time lowering the melting point [17] resulting in a better compatibility with sensitive substrates. The results of the simulations for five representative time steps are visualized in Fig. 2. It can be seen that both particle systems reach temperatures of 2000 °C or more in the ZnO layer by the end of the incident 30 ns laser pulse. While for stack B i.e. a ZnO particle layer without underlying buffer layer, the substrate is substantially heated up, for stack A the substrate stays comparatively cool under the SiO₂ particle layer. The color gradient in the buffer layer shows that heat transfer towards the substrate is hampered by the SiO₂ particle layer. It also becomes clear that the ZnO layer cools down

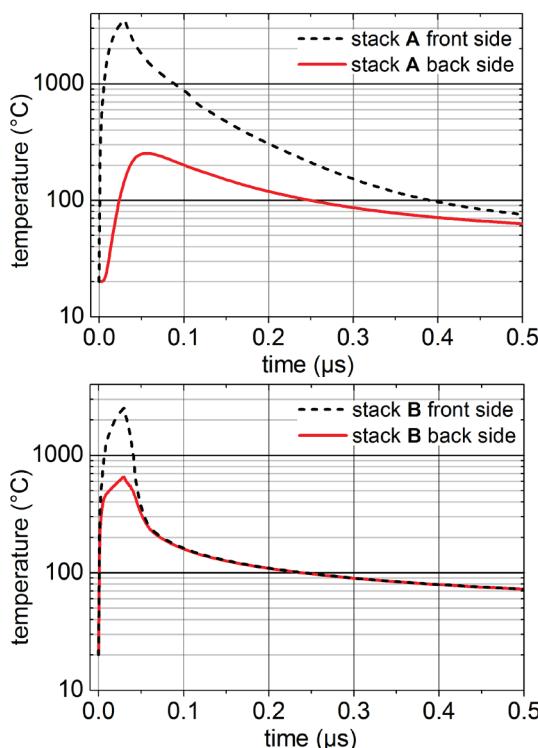


Fig. 3 Temperature course at the interface between air and the ZnO particle layer (“front side”) and at the interface between the particle layer stack and the substrate (“back side”) after irradiation with a 30 ns laser pulse beginning at $t = 0$ s with a fluence of 245 mJ/cm^2

more slowly for the layer stack with underlying buffer layer. A more quantitative representation of the outcome of the simulation is shown in fig. 3. It shows that particle stack B reaches a maximum temperature of about 2500 °C, a temperature well above the melting temperature of ZnO, while the maximum temperature of stack A is about 1000 °C higher which is a consequence of neglecting phase changes of the particle layer in the simulation. The cooling phase of the layer surface turns out to be significantly longer for stack A. The temperature right at the interface between the particle layer and the substrate is above 650 °C and shows steep gradients for stack B, while it stays below 250 °C and displays much lower gradients both in the heating and in the cooling phase for stack A. This outcome indicates that the same kind of buffer layer can be used to protect more sensitive substrates such as polymer foils which are required for the low-budget mass production of electron devices. It should be emphasized that the laser fluences needed for consolidation of the particle layer are reduced by applying a buffer layer and that a thicker buffer layer decreases the temperatures reached on the substrate side even more. The influence of the SiO_2 buffer layer thickness is shown in Fig. 4. The temperature course at the interface between the particle layer and the Si substrate was investigated in further simulations for different SiO_2 layer thicknesses (90 nm, 220 nm and 350 nm). It clearly shows that the maximum temperature reached at this interface is highly dependent on the buffer layer thickness. Thus, this layer should be tuned in thickness depending on the substrate material and maintained at a thickness high enough to guarantee substrate protection but at the same time low enough to not diminish the visual quality e.g. in display

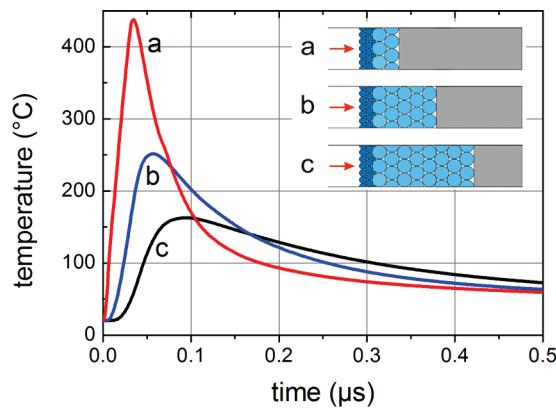


Fig. 4 Temperature course at the boundary between the particle layers and the Si substrate for different SiO_2 buffer thicknesses (**a** 93 nm, **b** 220 nm, **c** 350 nm). The inset shows schemes of the corresponding simulated layer stacks

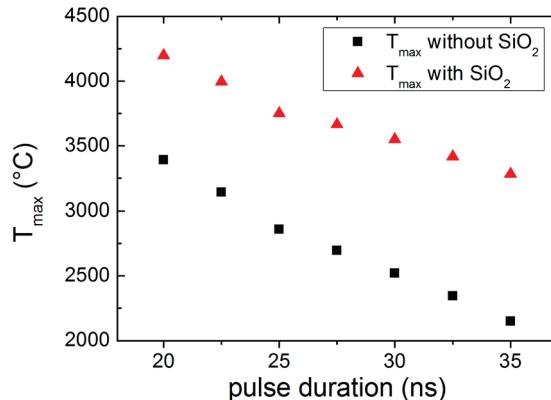


Fig. 5 Maximum temperatures reached at the boundary between air and the particle layer stack surface for particle layers with and without buffer layer for different pulse durations at a fluence of 245 mJ/cm^2 (stack A and B)

applications. The lasing gas mixture of excimer lasers such as the one we used ages over time due to the formation of chemical compounds of the highly reactive halogen gases. This leads to a decrease in pulse energy for a given excitation voltage as well as to a slight decrease of the pulse length on the order of 10 ns when the laser gas becomes older. This decrease in pulse energy can be compensated by increasing the excitation voltage. However, the pulse length remains shortened. While this fact is not such a big issue in the common fields of application of excimer lasers, such as lithography, it turned out to be a challenge when consolidated layers of unchanged quality shall be produced. Consequently, the influence of such a variation of the pulse length was also examined by the simulation model. The results for the maximum surface temperatures layers with and without a 220 nm thick SiO_2 buffer are shown in Fig 5. It can be seen that the maximum temperature reached at the layer surface varies by a few hundreds of degrees Celsius when the pulse length is changed by several nanoseconds. Thus the pulse length should be monitored during sample treatment. For layer production over several days, the process window will have to be adapted. Moreover, simulations with different gas atmospheres (Argon and Helium) were carried out. However, they showed that within the limits of the model i.e. only taking into account the thermal properties of the gases and neglecting other effects e.g. possible chemical reactions or changes in stoichiometry,

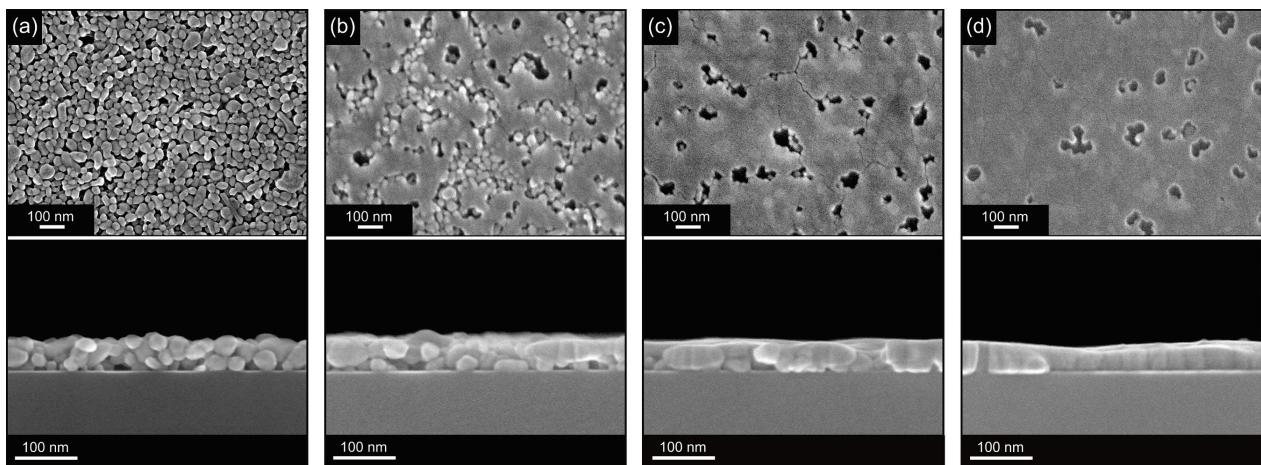


Fig. 6 Top-view and cross-sectional SEM images of different 40 nm thick ZnO particle layers on Si substrates. (a) untreated sample, (b) after irradiation by a single laser pulse at a fluence of 162 mJ/cm², (c) 187 mJ/cm², (d) 245 mJ/cm²

the type of gas atmosphere does only marginally change the maximum temperatures reached.

3. Experimental part

In our experiments, we used ZnO nanoparticles commercially available from Evonik Industries AG, with a broad particle size distribution ranging from 20 nm to 120 nm. An ethanol based dispersion was prepared from 20 wt% of these particles, while agglomeration of the particles was prevented by addition of 3 wt% 3,6,9-trioxadecanoic acid acting as a steric stabilizer. For the generation of a finer and more uniform particle size distribution, the suspensions were centrifuged resulting in an average particle size of about 25 nm. The preparation of the suspension is described in detail in [18]. For the generation of thermally insulating buffer layers we used the commercially available suspension 30V50 from Klebosol containing SiO₂ particles with an average particle diameter of about 50 nm. The samples were prepared on p-type silicon wafer substrates with a dimension of 20 mm x 20 mm and a thickness of

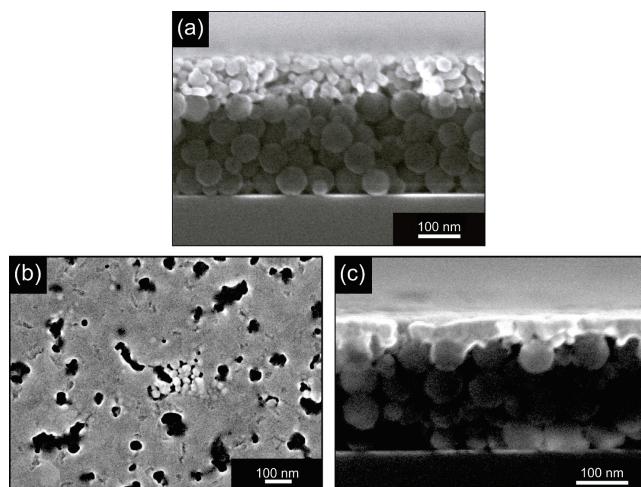


Fig. 7 (a) Cross-sectional SEM image of a ZnO particle layer (about 25 nm particle diameter) on a 220 nm SiO₂ particle layer on Si substrate. (b) Top-view and (c) cross-sectional SEM images of a layer stack consisting of a 40 nm ZnO particle layer on a 220 nm SiO₂ layer after irradiation with a laser pulse at a fluence of 60 mJ/cm²

1 mm. Prior to the deposition of the particle layers, the substrates were cleaned for 30 s in diluted HF (1 %) and then rinsed in deionized water. Layers of about 40 nm of ZnO particles or of about 220 nm SiO₂ particles followed by 40 nm ZnO particles were generated by spin coating the, suspensions on the substrates (Fig 6 a, Fig 7 a). SiO₂ particles were chosen for various reasons: they can be applied by the same technique as the ZnO particle layer, they do not significantly influence the optical appearance of the layer stack and feature a very low thermal conductivity. In order to thermally remove the stabilizing agents free of residues the samples were annealed at 400 °C for 30 min under ambient conditions. For the laser treatment of the particle layers, we used a Lambda Physik LPX 315i KrF excimer laser with a wavelength of 248 nm, a raw beam size of 30 mm x 15 mm, maximum pulse energy of 0.8 J and a pulse length of 30 ns while the temporal shape of the laser pulses is rectangular (we have to mention here that these values are measured for fresh laser gas. Obviously the maximum pulse energy decays over time, while the pulse length shortens slightly). The particle layers were irradiated by single pulses with fluences ranging from 60 mJ/cm² to 330 mJ/cm². While for fluences below 170 mJ/cm² the samples were irradiated by the raw beam which was limited to a size of 5 mm x 30 mm by an aperture, for fluences above 170 mJ/cm² a fused silica lens with a focal length of 500 mm was used to decrease the beam size. Fig. 6 shows the top and cross-sectional views of an untreated sample (a) as well as a series of samples treated with different increasing laser fluences (b) – (d). While laser fluences of 120 mJ/cm² and below (not shown here) do not seem to influence the morphology of the layer, fluences of 162 mJ/cm² and above lead to melting and significant densification of the particle layer. For a fluence of 245 mJ/cm², apart from some holes which might be a result of a not fully homogeneous initial distribution of particles, almost maximal density is reached while grain boundaries between the consolidated particles remain visible. A more detailed description of the shown layers and a comparison of the laser treated layers to conventionally annealed particle layers are given in [19]. An exemplary sample of ZnO particles on a SiO₂ buffer layer can be seen in Fig 7 (a). Experiments with such layer stacks show that the laser fluence needed to achieve a good con-

solidation of the ZnO layer are much lower in comparison to the samples without a SiO₂ buffer layer as indicated by the simulations. Fig 7 (b) and (c) show such a ZnO/SiO₂ layer stack after irradiation by a laser pulse at a fluence of only 60 mJ/cm². The surface and cross-section of the particle layer stack show that a very dense consolidation of the ZnO layer was achieved while the SiO₂ particles appear to be unaffected. The substantial decrease in the laser fluence needed for ZnO layer consolidation indicates that indeed the applied buffer layer can hamper the conduction of heat from the particle layer towards the substrate resulting in higher layer temperatures and as seen in the simulation part also increases the time needed for a cool down of the hot ZnO layer. A longer cool down time could contribute to enable diffusion processes and also to reduce thermal stress in the particle layer.

4. Conclusion

In conclusion, our experiments and simulations have shown that excimer laser irradiation is suitable for the consolidation of metal oxide nanoparticle layers. The laser fluences needed for consolidation of ZnO particle layers could be reduced greatly by adding a buffer layer of SiO₂ particles below the ZnO particle layer. These finding indicate that the transport of heat towards the substrate is significantly hindered by the buffer layer. Simulations of the system showed the same result. At the same time, the temperatures at the border between particle layer stack and substrate remain much lower. Thus, our investigations indicate that such a buffer layer could be used in the low-cost mass production of devices such as solar cells or displays on sensitive substrates like polymers despite the extremely high melting temperature of metal oxide nanoparticles such as ITO or ZnO while leaving the visual appearance of the layer unchanged. The performance of the buffer layer showed to be highly dependent on its thickness which has to be adapted to the substrate material.

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